MULTI-WAVELENGTH LIGHT SOURCE APPARATUS

Background of the Invention

Field of the Invention

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The present invention relates to a method and apparatus for easily implementing a multi-wavelength light source the frequency intervals of which are equal.

Description of the Related Art

10 The wavelength of signal light in a WDM (Wavelength Division Multiplexing) optical fiber communications system is stipulated to be arranged on a predetermined frequency grid by the ITU-T recommendations. Therefore, an absolute wavelength must be precisely controlled for an oscillation on this grid.

For a method preparing single wavelength lasers by a required number of channels, firstly, its monitoring/control becomes complicated. Secondly, it is inevitable to increase the size and the power consumption of an apparatus if the number of wavelengths, namely, the number of channels becomes large.

As a method for solving these problems, there is a method for splitting longitudinal mode components caused by modulation, and making the components into a multi-wavelength light source (see Non-Patent Document

1). The longitudinal mode is a spectrum component caused by modulation. If the spectrum of modulated light is viewed with a spectrum analyzer having a low resolution, it is shaped like a moderate mountain. However, if the spectrum is viewed with a spectrum analyzer having a high resolution, it is proved to be actually composed of many spectrum components having a narrow spectrum width. Each spectrum component having a narrow spectrum width, which configures such a spectrum of modulation light, is called a longitudinal mode component.

In Fig. 1, an optical pulse sequence from a pulse light source 10 that outputs an optical pulse sequence of a repetitive frequency f_0 Hz is input to a modulator array 11. In a wavelength demutiplexer 11-1, the longitudinal mode components of the optical pulse sequence are split and made into light beams having respective wavelengths. Then, the light beams are modulated by a modulator 11-3, and signals are put on the modulated light beams. Thereafter, these modulated light beams are coupled by a wavelength multiplexer 11-2, and transmitted.

Another characteristic of this method exists in a point that the number of channels can be increased by using spectrum broadening caused by nonlinear effects that occur within a nonlinear medium.

As conventional multi-wavelength light sources, techniques recited in Patent Documents 1 and 2 exist. With the technique recited in Patent Document 1, a longitudinal mode component obtained from modulated 1 light is demultiplexed, and made into a light source of each wavelength. Patent Document 2 discloses the technique with which an optical pulse sequence from a light source that generates an optical pulse sequence is passed through an optical fiber the dispersion of which is flattened to widen the width of a spectrum by nonlinear effects, and a longitudinal component is extracted from the widened spectrum.

[Patent Document 1]

Japanese Patent Application Publication No.

15 2001-264830

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[Patent Document 2]

Japanese Patent Application Publication No. 2002-236301

[Non-Patent Document]

20 IEEE Photonics Technology Letters, Vol. 9, No. 6, June 1997, pp. 818-820

Figs. 2A and 2B exemplify the configuration of an optical transmitting apparatus using another conventional multi-wavelength light source, and the shape of a spectrum.

In Fig. 2A, a pulse light source 15 outputs an optical pulse sequence of a repetitive frequency f_0 Hz, a spectrum expanding device 16 expands the spectrum of the optical pulse sequence, a modulation array then modulates each longitudinal mode component, and a gain equalizer 18 realizes the same power of each wavelength.

A plurality of single wavelength light sources can be created by extracting the longitudinal mode components of an optical spectrum with a narrowband filter as described above. With the conventional technique shown in Fig. 2A, however, the flatness of the spectrum of light after being broadened is poor, and the powers of respective signal wavelengths significantly vary.

A spectrum broadened by using a Gaussian pulse is exemplified in Fig. 2B. A horizontal axis represents a wavelength, whereas a vertical axis represents power on a linear scale. This spectrum is shaped like an envelope that connects the peaks of longitudinal mode components. Namely, this figure shows the spectrum viewed with a spectrum analyzer having a low resolution. As is known from this spectrum, powers vary by wavelength on the order of several times. Accordingly, it is difficult to create a practical multi-wavelength light source. To actually apply this multi-wavelength light source to a WDM communications system, a gain equalizer that

equalizes a power difference among channels must be incorporated. The gain equalizer matches the powers of wavelengths with that of a wavelength having the lowest power. Therefore, light of a wavelength originally having high power is attenuated, so that the loss of optical power increases, and also an optical signal to noise ratio is degraded.

Summary of the Invention

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An object of the present invention is to provide a method and an apparatus, which allow a plurality of single wavelength light sources to be obtained, and implements a multi-wavelength light source with which the output powers of respective wavelengths are made almost equal.

The multi-wavelength light source according to the present invention comprises: an optical pulse light source outputting an optical pulse sequence; an optical pulse shaping unit making the shape of an optical pulse output from the optical pulse light source into a super Gaussian pulse of the third order or higher; a spectrum expanding unit broadening the spectrum of an optical pulse sequence composed of shaped optical pulses; and a light splitting unit splitting the optical pulse sequence the spectrum of which is expanded into light

beams of respective frequencies.

According to the present invention, the shape of each pulse of an optical pulse sequence is made into a super Gaussian pulse of the third order or higher, whereby a spectrum obtained after being broadened has a good flatness, and a plurality of light beams having wavelengths the strengths of which are equal can be provided.

The multi-wavelength light source according to the
10 present invention can generate a flat spectrum. As a
result, a plurality of single wavelength light sources
can be provided without equalizing the optical power
of each wavelength.

15 Brief Description of the Drawings

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Fig. 1 exemplifies the configuration of a conventional multi-wavelength light source;

Figs. 2A and 2B exemplify the configuration of an optical transmitting apparatus using another conventional multi-wavelength light source, and the shape of a spectrum;

Figs. 3A and 3B show the basic configuration of a preferred embodiment according to the present invention;

Fig. 4 shows the basic configuration of a liquid

crystal spatial light modulator; and

Fig. 5 explains the definitions of symbols used in equations.

5 Description of the Preferred Embodiments

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A preferred embodiment according to the present invention overcomes the above described problems with the following means. Namely, the shape of each pulse in a pulse waveform sequence, which is obtained with modulation and whose spectrum is to be broadened, is made into a suitable shape such that the powers of respective wavelengths are made equal after the spectrum is broadened. Specifically, the pulse is made into a super Gaussian waveform of the third order or higher, its spectrum is expanded, then whereby a multi-wavelength optical source in which the powers of respective wavelengths are almost equal can be implemented.

The following description of the preferred embodiment mainly refers to a multi-wavelength light source using a pulse having a third-order super Gaussian waveform. However, the preferred embodiment is also applicable to a super Gaussian pulse of the third order or higher.

25 Figs. 3A and 3B show the basic configuration of

a preferred embodiment according to the present invention.

As shown in Fig. 3A, a third-order super Gaussian pulse is generated in a pulse light source 20, its spectrum 5 is expanded to be flat in a nonlinear medium of a spectrum expanding device 21, the light is split into respective wavelengths in a modulation array 22, and all of the wavelengths are again coupled after the data is modulated. Fig. 3B shows a light spectrum obtained by broadening 10 third-order super Gaussian pulses in the nonlinear medium. Also in this figure, the spectrum is shaped like an envelope that connects the peaks of longitudinal mode components, and does not show each longitudinal mode component. Since the spectrum shown in Fig. 3B has a 15 flatness in comparison with Fig. 2B, it is proved that using the third-order super Gaussian pulse (or a super Gaussian pulse of the third order or higher) allows a light source which makes powers almost equal to be obtained if the light is demultiplexed into respective 20 wavelengths.

Examples of a pulse light source include a semiconductor mode synchronous laser, a fiber ring laser, a semiconductor ring laser, a pulse light source using an electroabsorption modulator, etc. However, the pulse light source is not limited to these ones. Normally,

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pulses output from these pulse light sources are pulses of a Gaussian type or a sech type. Accordingly, in a light source, a pulse shaper making these pulses into a third order super Gaussian pulse waveform is required.

For this pulse light source for implementing a multi-wavelength light source, a short pulse on the order of several picoseconds is used to effectively cover a broad wavelength band. Since this pulse is faster than the operating speed of an electronic circuit, it cannot be followed with an electric operation method. However, a pulse shaper that shapes a wavelength in a frequency region as an optical signal left unchanged can be used. For example, a pulse shaper using a liquid crystal spatial light modulator (LC-SLM) exists. This shaper can generate a third-order super Gaussian pulse.

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Fig. 4 shows the basic configuration of the liquid crystal spatial light modulator.

Initially, pulses to be shaped are split into respective frequencies with a diffraction grating 30, and a focus is achieved on a Fourier plane with a convex lens 31. If a plurality of LC-SLMs 32, which can modulate the intensity and the phase of passing light, are arranged on the Fourier plane, the intensity and the pulse of the entire pulse band can be operated in a frequency region. After being operated, inverse Fourier transform

is performed for the pulses with a convex lens 33 and a diffraction grating 34, so that the pulses can be returned to a time domain. For the details of the principle of this waveform shaper, see the document "Opt. Lett. vol. 15, pp. 326-328, 1990". In principle, light beams having respective wavelengths of the light, for which Fourier transforming is performed in the diffraction grating 30, is passed through a liquid crystal spatial light modulator, suitable intensity and phase are given to each wavelength, and the inverse Fourier transform is performed for the light beams in the diffraction grating 33 to return as a waveform on a time axis, so that a desirable waveform can be obtained. Adjustments of the intensity and the phase for each wavelength are made by mathematically representing a preferable waveform, by performing the Fourier transform to calculate the intensity and the phase of each frequency or wavelength component, and by controlling the liquid crystal spatial light modulator based on the calculation.

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On the Fourier plane, a light resolution is determined by the characteristic of a diffraction grating, the beam diameter of light, and an incident angle to the diffraction grating. A larger value of the light resolution and a resolution determined by the width of crystal liquid spatial modulation becomes the resolution

of the Fourier transform. In the meantime, a bandwidth is determined uniquely by the focal distance and the diffraction angle of the convex lens.

Fig. 5 explains the definitions of symbols used 5 in equations.

Specifically, a resolution $(\delta\lambda)$, and a bandwidth $(\Delta\lambda)$, which are determined by the diffraction grating, etc., are represented by the following equations.

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$$\frac{\lambda}{\delta\lambda} = mN \frac{R}{\cos\theta}$$
$$\Delta\lambda = 2F\Delta\psi$$

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where λ indicates the wavelength of light, R indicates the beam diameter of the light, θ indicates the incident angle to the diffraction grating, N indicates the number of grooves (per unit length) of the diffraction grating, m indicates the order of the diffraction grating, which normally takes ± 1 , and F indicates the focal distance of the convex lens. As is known from Fig. 5, $\Psi(-)$ is an angle obtained by measuring the angle at the center of light reflected by the diffraction grating from a normal of the diffraction grating. The angle is orientated in a negative direction. $\delta\Psi$ indicates a deviation from the center of the reflected light as an angle. $\theta(-)$ indicates the incident angle

of incident light, which is measured from the normal of the diffraction grating, and the angle is orientated in the negative direction. $\Delta\Psi$ indicates a spreading angle of the light that is spectrum-decomposed and reflected by the diffraction grating.

If a diffraction grating having a large number of grooves is used, and if the incident angle is controlled to set $\cos\theta$ to a small value, a resolution of sub-nm can be achieved. In the meantime, if liquid crystal modulators the width of which is 0.1 mm are arranged by 128, the photofield of a bandwidth on the order of picoseconds can be operated.

Assuming that the electric field of a desired third-order super Gaussian pulse is $E_{m=3}\left(t\right)$, a transfer function $T\left(\omega\right)$ given by an LC-SLM is defined as follows.

$$T(\omega) = \frac{\widetilde{E}_{m=3}(\omega)}{\widetilde{E}_{0}(\omega)}$$

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where \sim indicates Fourier transform, and $E_0\left(t\right)$ indicates an incident pulse waveform.

As another method, there is a method performing Fourier transform with a planner lightwave circuit (PLC). In this case, diffraction the order of which is high can be implemented, so that Fourier transform can be made with a compact configuration. For the details, see

the reference "Y. Inoue, et al., IEEE Photonics Technology Letters, pp. 569-571, v. 11, no. 5, 1999". This document discloses an array waveguide grating having a capability similar to a diffraction grating.

In this preferred embodiment, the diffraction grating 30 may be a component that splits light. Similarly, the diffraction grating 34 may be a component that couples split light beams. Accordingly, these components are not limited to the diffraction gratings.

For example, a virtually imaged phased array (VIPA) element, which is disclosed by Japanese Patent Publication No. HEI09-043057, or the like, may be available as an alternative to a diffraction grating.

A spectrum spreading device is configured by a nonlinear medium having the third-order nonlinear effects. A spectrum is spread by the nonlinear effects within the medium. As the nonlinear medium, a highly nonlinear fiber whose nonlinear refractive index is improved by doping Ge, a holey fiber the nonlinearity of which is enhanced by reducing an effective core cross-sectional area with a plurality of holes made on the cross section of the fiber, or the like is effective. A spectrum spreading method is disclosed by Japanese Patent Publication No. 2002-77052.

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